Compensating a String PDE in the Actuation or Sensing Path of an Unstable ODE

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Abstract—How to control an unstable linear system with a long pure delay in the actuator path? This question was resolved using 'predictor' or 'finite spectrum assignment' designs in the 1970s. Here we address a more challenging question: How to control an unstable linear system with a wave PDE in the actuation path? Physically one can think of this problem as having to stabilize a system to whose input one has access through a string. The challenges of overcoming string/wave dynamics in the actuation path include their infinite dimension, finite propagation speed of the control signal, and the fact that all of their (infinitely many) eigenvalues are on the imaginary axis. In this paper we provide an explicit feedback law that compensates the wave PDE dynamics at the input of an LTI ODE and stabilizes the overall system. In addition, we prove robustness of the feedback to the error in a priori knowledge of the propagation speed in the wave PDE. Finally, we consider a dual problem where the wave PDE is in the sensing path and design an exponentially convergent observer.

I. INTRODUCTION

The 'Smith predictor' and its extensions developed since the 1970s [1], [3], [4], [5], [6], [7], [8], [13], [9], [14], [15], [16], [17], [18], [19], [21], [22], [25], [26], [27], [28], [29], [30], [31] are important tools in several application areas. They allow to compensate a pure delay of arbitrary length in either the actuation or sensing path of a linear system, even when the system is unstable. Several results in adaptive control for unknown ODE parameters have been published [2], [20]. Extensions to nonlinear systems are also beginning to emerge [10].

In [11] we presented a first attempt of compensating infinite-dimensional actuator dynamics of more complex type than pure delay. We presented a design for diffusiondominated PDE dynamics (such as the heat equation). While these dynamics do not have a finite speed of propagation, they are 'low-pass' and 'phase-lag' to the extreme, as they have infinitely many (stable) poles and no zeros.

In this paper we tackle a problem from a different class of PDE dynamics in the actuation or sensing path—the wave/string equation. The wave equation is challenging due to the fact that all of its (infinitely many) eigenvalues are on the imaginary axis, and due to the fact that it has a finite (limited) speed of propagation (large control doesn't help).

The problem studied here is more challenging than in [11] due to another difficulty—the PDE system is second order in time, which means that the state is 'doubly infinite dimensional' (distributed displacement and distributed velocity). This is not so much or a problem dimensionally, as it is a problem in constructing the state transformations for compensating the PDE dynamics. One has to deal with the coupling of two infinite-dimensional states.

As in [13] for delay-ODE cascades, and in [11] for heat-PDE-ODE cascades, we design feedback laws that are given by explicit formulae. We start in Section II with an actuator compensation design with full state feedback. In Section III we approach the question of robustness of our infinite-dimensional feedback law with respect to small uncertainty in the wave propagation speed and provide an affirmative answer. Finally, in Section IV we develop a dual of our actuator dynamics compensator and design an infinite-dimensional observer which compensates the wave PDE dynamics of the sensor.

II. STABILIZATION WITH FULL-STATE FEEDBACK

We consider the cascade of a wave (string) equation and an LTI finite-dimensional system given by

$$\dot{X}(t) = AX(t) + Bu(0,t) \tag{1}$$

$$u_{tt}(x,t) = u_{xx}(x,t) \tag{2}$$

$$u_x(0,t) = 0 \tag{3}$$

$$u_x(D,t) = U(t), \qquad (4)$$

where $X \in \mathbb{R}^n$ is the ODE state, *U* is the scalar input to the entire system, and u(x,t) is the state of the PDE dynamics of the actuator governed by a wave equation. The cascade system is depicted in Figure 1.

The length of the PDE domain, D, is arbitrary. Thus, we take the wave propagation speed to be unity without loss of generality. We assume that the pair (A,B) is stabilizable and take K to be a known vector such that A + BK is Hurwitz.

We recall from [13] that, if (2), (3) are replaced by the delay/transport equation,

$$u_t(x,t) = u_x(x,t), \qquad (5)$$

then the predictor-based control law

$$U(t) = K \left[e^{AD} X(t) + \int_0^D e^{A(D-y)} Bu(y,t) dy \right]$$
(6)

achieves perfect compensation of the actuator delay and achieves exponential stability at $u \equiv 0, X = 0$. When the pure delay actuator dynamics are replaced by the wave equation dynamics, a much more involved feedback law is needed.

We seek an invertible transformation $(X, u, u_t) \mapsto (X, v, v_t)$ that converts (1)–(3) into

$$\dot{X}(t) = (A + BK)X(t) + Bv(0, t)$$
(7)

$$v_{tt}(x,t) = v_{xx}(x,t) \tag{8}$$

$$v_x(0,t) = 0 \tag{9}$$

$$U(t) \xrightarrow{u_{\chi}(D,t)} \underbrace{\begin{array}{c}u(x,t)\\wave eqn. PDE\\(actuator)\end{array}} \underbrace{u(0,t)}_{(plant)} \underbrace{\begin{array}{c}V(t)\\odd \\(plant)\end{array}}$$

Fig. 1. The cascade of the wave equation PDE dynamics of the actuator with the ODE dynamics of the plant.

and then another transformation $(X, v, v_t) \mapsto (X, w, w_t)$ that converts (7)–(9) into

$$\dot{X}(t) = (A + BK)X(t) + Bw(0,t)$$
(10)

$$w_{tt}(x,t) = w_{xx}(x,t) \tag{11}$$

$$w_x(0,t) = c_0 w(0,t), \qquad c_0 > 0.$$
 (12)

We also seek a feedback law that achieves

$$w_x(D,t) = -c_1 w_t(D,t), \qquad c_1 > 0.$$
 (13)

The system (10)–(13) is exponentially stable, as we shall see. With the invertibility of the composite transformation $(X, u, u_t) \mapsto (X, w, w_t)$, we will achieve exponential stability of the closed-loop system in the original variables (X, u, u_t) .

We postulate the transformation $(X, u, u_t) \mapsto (X, v, v_t)$ in the form

$$v(x,t) = u(x,t) - \int_0^x k(x,y)u(y,t)dx - \int_0^x l(x,y)u_t(y,t)dy - \gamma(x)X(t),$$
(14)

where the kernel functions k(x,y), l(x,y), and $\gamma(x)$ are to be found. By matching the systems (1)–(3) and (7)–(9), a lengthy but straightforward calculation leads to the following conditions on the kernels:

$$\gamma''(x) = \gamma(x)A^2 \tag{15}$$

$$\gamma(0) = K \tag{16}$$

$$\gamma'(0) = 0 \tag{17}$$

$$l_{xx}(x,y) = l_{yy}(x,y)$$
 (18)

$$l(x,x) = 0 (19)$$

$$I_{y}(x,0) = -\gamma(x)B \tag{20}$$

$$k_{xx}(x,y) = k_{yy}(x,y)$$
 (21)
 $k(x,y) = 0$ (22)

$$k(x,x) = 0$$
(22)

$$k_y(x,0) = -\gamma(x)AB. \qquad (23)$$

These differential equations can be solved explicitly. The solutions are

$$\gamma(x) = KM(x) \tag{24}$$

$$l(x,y) = m(x-y)$$
(25)
$$k(x,y) = \mu(x-y)$$
(26)

$$k(x,y) = \mu(x-y)$$

$$\begin{bmatrix} 0 & A^2 \end{bmatrix}$$
(26)

$$M(x) = \begin{bmatrix} I & 0 \end{bmatrix} e^{\begin{bmatrix} I & 0 \end{bmatrix}^{x}} \begin{bmatrix} I \\ 0 \end{bmatrix}$$
(27)

$$m(s) = \int_0^s \gamma(\xi) B d\xi$$
 (28)

$$\mu(s) = \int_0^s \gamma(\xi) ABd\xi.$$
 (29)

Thus the transformation $(X, u, u_t) \mapsto (X, v, v_t)$ is defined as

$$v(x,t) = u(x,t) - \int_0^x \mu(x-y)u(y,t)dx - \int_0^x m(x-y)u_t(y,t)dy - \gamma(x)X(t)$$
(30)

$$v_t(x,t) = u_t(x,t) - KBu(x,t) - \int_0^x \mu(x-y)u_t(y,t)dx - \int_0^x m''(x-y)u(y,t)dy - \gamma(x)AX(t).$$
(31)

With similar derivations, one can show that the inverse of the transformation $(X, u, u_t) \mapsto (X, v, v_t)$ is defined as

$$u(x,t) = v(x,t) - \int_0^x \sigma(x-y)v(y,t)dx - \int_0^x n(x-y)v_t(y,t)dy - \rho(x)X(t)$$
(32)

$$u_{t}(x,t) = v_{t}(x,t) + KBv(x,t) - \int_{0}^{x} \sigma(x-y)v_{t}(y,t)dx - \int_{0}^{x} n''(x-y)v(y,t)dy - \rho(x)AX(t), \quad (33)$$

where

$$\rho(x) = -KN(x) \qquad (34)$$
$$\begin{bmatrix} 0 & (A+BK)^2 \end{bmatrix}$$

$$N(x) = \begin{bmatrix} I & 0 \end{bmatrix} e^{\begin{bmatrix} 0 & (II + DII) \\ I & 0 \end{bmatrix}^{x}} \begin{bmatrix} I \\ 0 \end{bmatrix}$$
(35)

$$n(s) = \int_0^s \rho(\xi) B d\xi \tag{36}$$

$$\sigma(s) = \int_0^s \rho(\xi) ABd\xi.$$
(37)

The transformation $(X, v, v_t) \mapsto (X, w, w_t)$ is simpler and given by

$$w(x,t) = v(x,t) + c_0 \int_0^x v(y,t) dy$$
 (38)

$$w_t(x,t) = v_t(x,t) + c_0 \int_0^x v_t(y,t) dy,$$
 (39)

whereas its inverse is

$$v(x,t) = w(x,t) - c_0 \int_0^x e^{-c_0(x-y)} w(y,t) dy$$
 (40)

$$v_t(x,t) = w_t(x,t) - c_0 \int_0^x e^{-c_0(x-y)} w_t(y,t) dy.$$
 (41)

4098

The composite transformation $(X, u, u_t) \mapsto (X, w, w_t)$ is

$$w(x,t) = u(x,t) + \int_0^x \left(c_0 - \mu(x-y) - c_0 \int_0^{x-y} \mu(\xi) d\xi \right) u(y,t) dy - \int_0^x \left(m(x-y) + c_0 \int_0^{x-y} m(\xi) d\xi \right) u_t(y,t) dy - \left(\gamma(x) + c_0 \int_0^x \gamma(\xi) d\xi \right) X(t)$$

$$w_t(x,t) = u_t(x,t) - KBu(x,t)$$
(42)

$$v_{t}(x,t) = u_{t}(x,t) - KBu(x,t) - \int_{0}^{x} \left(c_{0}m'(x-y) + m''(x-y) \right) u(y,t) dy + \int_{0}^{x} \left(c_{0} - \mu(x-y) - c_{0} \int_{0}^{x-y} \mu(\xi) d\xi \right) u_{t}(y,t) dy - \left(\gamma(x) + c_{0} \int_{0}^{x} \gamma(\xi) d\xi \right) AX(t)$$
(43)

and its inverse is

$$u(x,t) = w(x,t) - \int_{0}^{x} \left(c_{0}e^{-c_{0}(x-y)} + \sigma(x-y) \right) - c_{0} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)} \sigma(\xi) d\xi w(y,t) dy - \int_{0}^{x} \left(n(x-y) - c_{0} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)} n(\xi) d\xi \right) \times w_{t}(y,t) dy - \rho(x)X(t)$$
(44)
$$u_{t}(x,t) = w_{t}(x,t) + KBw(x,t) - \int_{0}^{x} \left(n''(x-y) - c_{0}n'(x-y) + c_{0}^{2}n(x-y) \right) + c_{0}^{3} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)} n(\xi) d\xi w(y,t) dy - \int_{0}^{x} \left(c_{0}e^{-c_{0}(x-y)} + \sigma(x-y) \right) - c_{0} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)} \sigma(\xi) d\xi w_{t}(y,t) dy - \rho(x)AX(t).$$
(45)

Next, we design a controller that satisfies the boundary condition (13). First, from (42) we get

$$w_{x}(x,t) = u_{x}(x,t) + c_{0}u(x,t) - \int_{0}^{x} \left(\mu'(x-y) + c_{0}\mu(x-y) \right) u(y,t) dy - \int_{0}^{x} \left(m'(x-y) + c_{0}m(x-y) \right) u_{t}(y,t) dy - \left(\gamma'(x) + c_{0}\gamma(x) \right) X(t) .$$
(46)

Then, the control law is

$$U(t) = (-c_0 + c_1 KB)u(D, t) - c_1 u_t(D, t) + \int_0^D p(D - y)u(y, t)dy + \int_0^D q(D - y)u_t(y, t)dy + \pi(D)X(t)$$
(47)

where

$$p(s) = \mu'(s) + c_0\mu(s) + c_1\left(m''(s) + c_0m'(s)\right)$$
(48)

$$q(s) = m'(s) + c_0 m(s) + c_1 \left(\mu(s) + c_0 \int_0^s \mu(\xi) d\xi - c_0 \right)$$
(49)

$$\pi(x) = \gamma'(x) + \gamma(x)(c_0 I + c_1 A) + c_1 c_0 \int_0^x \gamma(\xi) d\xi A.$$
 (50)

Next we state a new controller that compensates the wave PDE actuator dynamics and prove exponential stability of the resulting closed-loop system.

Theorem 1: (Stabilization) Consider a closed-loop system consisting of the plant (1)–(4) and the control law (47). For any initial condition such that $u(\cdot,0) \in H^1$ and $u_t(\cdot,0) \in L^2$, the closed-loop system has a unique solution $(X(t), u(\cdot,t), u_t(\cdot,t)) \in C([0,\infty), \mathbb{R}^n \times H^1(0,D) \times L^2(0,D))$ and is exponentially stable in the sense of the norm

$$\left(|X(t)|^2 + u(0,t)^2 + \int_0^D u_x(x,t)^2 dx + \int_0^D u_t(x,t)^2 dx\right)^{1/2}.$$
(51)

Moreover, if the initial condition $(u(\cdot,0), u_t(\cdot,0))$ is compatible with the control law (47) and belongs to $H^2(0,D) \times H^1(0,D)$, then $(X(t), u(\cdot,t), u_t(\cdot,t)) \in C^1([0,\infty), \mathbb{R}^n \times H^1(0,D) \times L^2(0,D))$ is the classical solution of the closed-loop system.

Proof: We will use the system norms

$$\Omega(t) = u(0,t)^{2} + ||u_{x}(t)||^{2} + ||u_{t}(t)||^{2} + |X(t)|^{2}$$
(52)
$$\Xi(t) = w(0,t)^{2} + ||w_{x}(t)||^{2} + ||w_{t}(t)||^{2} + |X(t)|^{2},$$
(53)

where $||u(t)||^2$ is a compact notation for $\int_0^D u(x,t)^2 dx$. In addition, we employ a Lyapunov function

$$V(t) = X(t)^T P X(t) + a E(t),$$
 (54)

where the matrix $P = P^T > 0$ is the solution to the Lyapunov equation $P(A+BK) + (A+BK)^T P = -Q$ for some $Q = Q^T > 0$, the parameter a > 0 is to be chosen later, and the function E(t) is defined by

$$E(t) = \frac{1}{2} \left(c_0 w(0,t)^2 + \|w_x(t)\|^2 + \|w_t(t)\|^2 \right) \\ + \delta \int_0^D (1+x) w_x(y,t) w_t(y,t) dy,$$
(55)

By using (43), (46), (45), and

$$u_{x}(x,t) = w_{x}(x,t) - c_{0}w(x,t) - \int_{0}^{x} \left(-c_{0}^{2}e^{-c_{0}(x-y)} + \sigma'(x-y) - c_{0}\sigma(x-y) \right) + c_{0}^{2} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)}\sigma(\xi)d\xi w(y,t)dy - \int_{0}^{x} \left(n'(x-y) - c_{0}n(x-y) \right) + c_{0}^{2} \int_{0}^{x-y} e^{-c_{0}(x-y-\xi)}n(\xi)d\xi w_{t}(y,t)dy - \rho'(x)X(t)$$
(56)

$$u(0,t) = w(0,t) + KX(t), \qquad (57)$$

and by using Poincare's inequality, for sufficiently small δ it is possible to show that there exist positive constants $\theta_1, \theta_2, \theta_3, \theta_4$ such that

$$\theta_1 \Xi \leq \Omega \leq \theta_2 \Xi$$
 (58)

$$\theta_3 \Xi \leq V \leq \theta_4 \Xi. \tag{59}$$

Furthermore, it is readily shown that

$$\dot{E}(t) = -\left(c_1 - \delta \frac{1+D}{2} \left(1+c_1^2\right)\right) w_t(D,t)^2 -\frac{\delta}{2} \left(w_t(0,t)^2 + c_0^2 w(0,t)^2\right) -\frac{\delta}{2} \left(\|w_x(t)\|^2 + \|w_t(t)\|^2\right).$$
(60)

Then, by choosing

$$a \ge \frac{8|PB|^2}{\delta c_0^2 \lambda_{\min}(Q)},\tag{61}$$

we get

$$\dot{V} \le -\eta V \tag{62}$$

for some sufficiently small positive η . From (58), (59), (62), it follows that

$$\Omega(t) \le \frac{\theta_2 \theta_4}{\theta_1 \theta_3} \Omega(0) \mathrm{e}^{-\eta t} \,. \tag{63}$$

The rest of the argument is almost identical to [12].

III. ROBUSTNESS TO UNCERTAINTY IN THE WAVE PROPAGATION SPEED

We now study robustness of the feedback law (47) to a small perturbation of the propagation speed in the actuator dynamics, i.e., we study stability robustness of the closedloop system

$$\dot{X}(t) = AX(t) + Bu(0,t) \tag{64}$$

$$u_{tt}(x,t) = (1+\varepsilon)u_{xx}(x,t) \tag{65}$$

$$u_x(0,t) = 0 \tag{66}$$

$$u_{x}(D,t) = (-c_{0} + c_{1}KB)u(D,t) - c_{1}u_{t}(D,t) + \int_{0}^{D} p(D-y)u(y,t)dy + \int_{0}^{D} q(D-y)u_{t}(y,t)dy + \pi(D)X(t)$$
(67)

to the perturbation parameter ε , which we allow to be either positive or negative but small.

With a very long calculation, we arrive at the representation of the system (64)–(67) in the *w*-variable:

$$\dot{X}(t) = (A + BK)X(t) + Bw(0,t)$$
 (68)

$$w_{tt}(x,t) = (1+\varepsilon)w_{xx}(x,t) + \varepsilon\Pi(x,t)$$
(69)

$$w_x(0,t) = c_0 w(0,t) \tag{70}$$

$$w_x(D,t) = -c_1 w_t(D,t),$$
(71)

where

$$\Pi(x,t) = \left(\gamma(x) + c_0 \int_0^x \gamma(\xi) d\xi\right)$$

$$\times \left((A + BK)X(t) + Bw(0,t)\right)$$

$$+ KBw_t(x,t) + (KB)^2w(x,t)$$

$$+ \int_0^x g(x-y)w(y,t)dy + \int_0^x h(x-y)w_t(y,t)dy$$

$$- \left(KB\rho(x) + \int_0^x \omega(x-\xi)\rho(\xi)d\xi\right)AX(t) \quad (72)$$

and where

$$\omega(x) = m''(x) + c_0 m'(x)$$
 (73)

$$g(x) = KB\phi(x) + KB\omega(x) + \int_0^x \omega(x-y)\phi(y)dy$$
(74)

$$h(x) = KB\psi(x) + \omega(x) + \int_0^x \omega(x - y)\psi(y)dy$$
(75)

$$\phi(x) = -n''(x) + c_0 n'(x) - c_0^2 n(x) + c_0^3 \int_0^x e^{-c_0(x-y)} n(y) dy$$
(76)

$$\Psi(x) = -c_0 e^{-c_0 x} - \sigma(x) + c_0 \int_0^x e^{-c_0(x-y)} \sigma(y) dy.$$
(77)

The state perturbation $\Pi(x,t)$ is very complicated in appearance but $\int_0^D \Pi(x,t)^2 dx$ can be bounded in terms of $\Xi(t)$ as defined in (53), and hence also in terms of V(t) as defined in (54). Consequently, the same kind of Lyapunov analysis can be conducted as in the proof of Theorem 1, dominating the effect of $\Pi(x,t)$ for small ε , to prove the following robustness result.

Theorem 2: (Robustness to Small Error in Wave Propagation Speed) Consider the closed-loop system (64)–(67). There exists a sufficiently small $\varepsilon^* > 0$ such that for all $\varepsilon \in (-\varepsilon^*, \varepsilon^*)$ the closed-loop system is exponentially stable in the same sense as in Theorem 1.

IV. OBSERVER DESIGN

Consider the LTI ODE system in cascade with a wave PDE in the sensing path (as depicted in Figure 2),

$$Y(t) = u(0,t)$$
 (78)

$$u_{tt}(x,t) = u_{xx}(x,t)$$
 (79)

$$u_x(0,t) = 0$$
 (80)

$$u(D,t) = CX(t) \tag{81}$$

$$\dot{X}(t) = AX(t) + BU(t).$$
(82)

We recall from [13] that, if (79), (80) are replaced by the delay/transport equation, $u_t(x,t) = u_x(x,t)$, then the predictor-based observer

$$\hat{u}_t(x,t) = \hat{u}_x(x,t) + C e^{Ax} L(Y(t) - \hat{u}(0,t))$$
(83)

$$\hat{u}(D,t) = C\hat{X}(t) \tag{84}$$

$$\hat{X}(t) = A\hat{X}(t) + BU(t) + e^{AD}L(Y(t) - \hat{u}(0,t))$$
(85)

achieves perfect compensation of the observer delay and achieves exponential stability at $u - \hat{u} \equiv 0, X - \hat{X} = 0$.

$$U(t) \longrightarrow \begin{bmatrix} X(t) & u(x,t) \\ ODE & u(D,t) \\ (plant) & CX(t) \end{bmatrix} \xrightarrow{u(x,t)} u(0,t) \xrightarrow{u(0,t)} Y(t)$$

Fig. 2. The cascade of the ODE dynamics of the plant with the wave equation PDE dynamics of the sensor.

We are seeking an observer of the form

$$\hat{u}_t(x,t) = \hat{u}_{xx}(x,t) + \alpha(x) \left(Y(t) - \hat{u}(0,t) \right) + \beta(x) \left(\dot{Y}(t) - \hat{u}_t(0,t) \right)$$
(86)

$$\hat{u}_x(0,t) = -a(Y(t) - \hat{u}(0,t)) - b(\dot{Y}(t) - \hat{u}_t(0,t))$$
(87)

$$\hat{u}(D,t) = C\hat{X}(t) \tag{88}$$

$$\hat{X}(t) = A\hat{X}(t) + BU(t) + \Lambda(Y(t) - \hat{u}(0, t)), \quad (89)$$

where the functions $\alpha(x)$, $\beta(x)$, the scalars *a*, *b*, and the vector Λ are to be determined, to achieve exponential stability of the observer error system

$$\tilde{u}_t(x,t) = \tilde{u}_{xx}(x,t) - \alpha(x)\tilde{u}(0,t) - \beta(x)\tilde{u}_t(0,t) \quad (90)$$

$$\tilde{u}_x(0,t) = a\tilde{u}(0,t) + b\tilde{u}_t(0,t)$$
(91)

$$\tilde{u}(D,t) = C\tilde{X}(t) \tag{92}$$

$$\dot{\tilde{X}}(t) = A\tilde{X}(t) - \Lambda \tilde{u}(0,t), \qquad (93)$$

where

$$\tilde{u}(x,t) = u(x,t) - \hat{u}(x,t)$$
(94)

$$\dot{X}(t) = X(t) - \dot{X}(t).$$
 (95)

We consider the transformation

$$\tilde{w}(x) = \tilde{u}(x) - \Gamma(x)\tilde{X}$$
(96)

and try to find $\Gamma(x)$, along with $\alpha(x), \beta(x), a, b, \Lambda$, that convert (90)–(93) into the exponentially stable system

$$\tilde{w}_{tt}(x,t) = \tilde{w}_{xx}(x,t) \tag{97}$$

$$\tilde{w}_x(0,t) = c_0 \tilde{w}_t(0,t) \tag{98}$$

$$\tilde{w}(D,t) = 0 \tag{99}$$

$$\tilde{X}(t) = (A - \Lambda \Gamma(0))\tilde{X} - \Lambda \tilde{w}(0, t), \qquad (100)$$

where $c_0 > 0$ and $A - \Lambda \Gamma(0)$ is a Hurwitz matrix.

By matching the systems (90)–(93) and (97)–(100), we obtain the conditions

$$\Gamma''(x) = \Gamma(x)A^2 \tag{101}$$

$$\Gamma'(0) = c_0 \Gamma(0) A \tag{102}$$

$$\Gamma(D) = C, \tag{103}$$

as well as

$$\alpha(x) = \Gamma(x)A\Lambda \tag{104}$$

$$\beta(x) = \Gamma(x)\Lambda \tag{105}$$

$$a = c_0 \Gamma(0) \Lambda \tag{106}$$

$$b = c_0.$$
 (107)

Solving the linear ODE two-point-boundary-value problem (101)-(103), we obtain

$$\Gamma(x) = \Gamma(0)G(x), \qquad (108)$$

where

$$\Gamma(0) = CG(D)^{-1} \tag{109}$$

$$G(x) = \begin{bmatrix} I & c_0 A \end{bmatrix} e^{\begin{bmatrix} 0 & A \\ I & 0 \end{bmatrix}^x} \begin{bmatrix} I \\ 0 \end{bmatrix}. \quad (110)$$

Thus, we have determined all the quantities needed to implement the observer (86)–(89) except Λ , which needs to be chosen so that the matrix $A - \Lambda\Gamma(0)$ is Hurwitz. We pick Λ as

$$\Lambda = G(D)L, \tag{111}$$

where *L* is chosen so that the matrix A - LC is Hurwitz. Since *A* and G(D) commute, using G(D) as a similarity transformation for the matrix $A - \Lambda\Gamma(0) = A - G(D)LCG(D)^{-1}$, we get that the matrices A - LC and $A - \Lambda\Gamma(0)$ have the same eigenvalues, so the latter matrix is Hurwitz.

So the system (97)-(100) is a cascade of a wave equation (97)-(99), which is exponentially stable due to the 'damping' boundary condition (98), and of the exponentially stable ODE (100). So, the entire observer error system is exponentially stable.

Theorem 3: (Observer Design and Convergence) Assume that the matrix G(D) is non-singular. The observer (86)–(89), with gains defined through (104)–(111), guarantees that \hat{X} , \hat{u} exponentially converge to X, u, i.e., more precisely, that the observer error system is exponentially stable in the sense of the norm

$$(|X(t) - \hat{X}(t)|^{2} + \int_{0}^{D} (u_{x}(x,t) - \hat{u}_{x}(x,t))^{2} dx + \int_{0}^{D} (u_{t}(x,t) - \hat{u}_{t}(x,t))^{2} dx \right)^{1/2}.$$
 (112)

Proof: Very similar to the proof of Theorem 1, with a Lyapunov function

$$V(t) = \tilde{X}(t)^{T} M(D)^{-T} P M(D)^{-1} \tilde{X}(t) + aE(t), \qquad (113)$$

where $P = P^T > 0$ is the solution to the Lyapunov equation $P(A - LC) + (A - LC)^T P = -Q$ for some $Q = Q^T > 0$, and with

$$E(t) = \frac{1}{2} \left(\|\tilde{w}_{x}(t)\|^{2} + \|\tilde{w}_{t}(t)\|^{2} \right) \\ + \delta \int_{0}^{D} (-1 - D + x) \tilde{w}_{x}(y, t) \tilde{w}_{t}(y, t) dy, (114)$$

4101

The system norms are simpler,

$$\Omega(t) = \|\tilde{u}_x(t)\|^2 + \|\tilde{u}_t(t)\|^2 + |\tilde{X}(t)|^2$$
(115)

$$\Xi(t) = \|\tilde{w}_x(t)\|^2 + \|\tilde{w}_t(t)\|^2 + |\tilde{X}(t)|^2, \quad (116)$$

and the system transformations are much simpler,

$$\widetilde{w}_{x}(x,t) = \widetilde{u}_{x}(x,t) - \Gamma'(x)\widetilde{X}(t)$$
(117)
$$\widetilde{w}_{t}(x,t) = \widetilde{u}_{t}(x,t) - \Gamma(x)A\widetilde{X}(t) + \Gamma(x)\Lambda\widetilde{u}(0,t)$$

$$\tilde{u}_t(x,t) = \tilde{w}_t(x,t) + \Gamma(x)(A - \Lambda \Gamma(0))\tilde{X}(t) -\Gamma(x)\Lambda \tilde{w}(0,t).$$
(118)

One obtains the inequalities (58), (59) with the help of Agmon's inequality, or, with the help of Poincare's inequality and the alternative representation of the state transformation,

$$\begin{split} \tilde{w}_t(x,t) &= \tilde{u}_t(x,t) + \Gamma(x)\Lambda \tilde{u}(x,t) - \Gamma(x)\Lambda \int_0^D \tilde{u}_x(y,t)dy \\ &-\Gamma(x)A\tilde{X}(t) \\ \tilde{u}_t(x,t) &= \tilde{w}_t(x,t) - \Gamma(x)\Lambda \tilde{w}(x,t) + \Gamma(x)\Lambda \int_0^d \tilde{w}_x(y,t)dy \\ &+\Gamma(x)(A - \Lambda \Gamma(0))\tilde{X}(t) \,. \end{split}$$
(119)

Then, one obtains (63), which completes the proof.

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