Asymptotic stability of stationary states of a stochastic neural field in the form of a PDE

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Workshop on nonlocal and nonlinear PDEs Trondheim 24.05.2023



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- Proposed as a model to study the effects of noise in grid cells.<sup>1</sup>
- The brain is noisy.<sup>3</sup>
- Varying the noise level can shift the behaviour of the network from one state to another.
- Existence of bifurcation branches gives possible stationary states, their stability tells us which ones the model could settle into.



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Assume  ${\cal B}$  constant. Stationary states satisfy

$$\sigma \partial_s \rho(x,s) = -\left(s - \Phi_{\bar{\rho}}\right) \rho(x,s).$$

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Hence, they solve

$$\rho = \frac{1}{Z} e^{-\frac{(s-\Phi_{\bar{\rho}})^2}{2\sigma}}.$$

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With conservation of unit mass we have

$$Z = \int_0^{+\infty} e^{-\frac{\left(s - \Phi_{\bar{\rho}}\right)^2}{2\sigma}} ds.$$

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If B > 0 and  $\int_{\mathbb{T}^d} W(x) dx < 0$ , and  $\Phi' \ge 0$ , the homogeneous in space stationary states are unique.<sup>1</sup>

<sup>&</sup>lt;sup>1</sup>Carrillo, Holden, S., JOMB 2022.

## Noise-induced apparition of patterns<sup>2</sup>

Let  $\rho_{\infty}$  denote the spatially homogeneous stationary state, and  $\frac{1}{\sigma_0} > \frac{2|W_0|^2}{\pi B^2}$ . Assume  $\Phi'' > -C_{\sigma_0}$  and that  $\exists k^* \in \mathbb{N}^d$  s.t.

$$\hat{W}_{k^*} = \frac{\sigma_0}{\Phi_0' M_\infty},$$

with

$$\Phi_0' = \Phi' \big( W_0 \bar{\rho}_\infty + B \big), \quad M_\infty = \int_0^{+\infty} (s - \bar{\rho}_\infty)^2 \rho_\infty(s) ds,$$

Then there exists spatially patterned bifurcation branches emanating from  $(\rho_{\infty}^{\sigma_0}, \sigma_0)$ .



<sup>2</sup>Carrillo, Roux, S., Physica D 2023.

## Possible patterns at the first three bifurcation points



# Numerical bifurcation diagram<sup>2</sup>



<sup>2</sup>Carrillo, Roux, S., Physica D 2023.

## Stability of stationary states

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But we can say something about the spatially homogeneous stationary states, and *maybe* other states.

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- 1. Linearise the PDE around  $\rho_\infty.$
- 2. Carefully combine estimates for the time derivative of the relative entropy of each Fourier mode  $(\hat{\rho} \hat{\rho}_{\infty})_k(s,t)$  of the perturbation  $(\rho \rho_{\infty})(x, s, t)$ , and the square of the Fourier modes of the perturbation of the mean,  $(\hat{\rho} \hat{\rho}_{\infty})_k$ .

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Resulted in exponential decay of the quantities

$$\int_0^{+\infty} \left( \frac{(\hat{\rho} - \hat{\rho}_\infty)_k}{\rho_\infty} \right)^2 \rho_\infty ds - \frac{\Phi'_0}{\sigma} \hat{W}_k ((\hat{\rho} - \hat{\rho}_\infty)_k)^2, \qquad k \in \mathbb{Z}^d,$$

which where shown to be positive under the condition

$$\hat{W}_k < \frac{\sigma}{\Phi'_0 M_\infty}.$$

<sup>1</sup>Carrillo, Holden, S., JOMB 2022

## An optimal condition for linear stability<sup>1</sup>

The spatially homogeneous stationary state  $\rho_{\infty}$ , which depends on  $\sigma$ , is *linearly exponentially stable* in the  $L^1$ -norm of the relative entropy,

$$\int_{\mathbb{T}^d} \int_0^{+\infty} \left( \frac{\rho(x, s, t) - \rho_{\infty}(s)}{\rho_{\infty}(s)} \right)^2 \rho_{\infty}(s) ds dx,$$

when all the unnormalised Fourier modes of the component-wise symmetric  $\boldsymbol{W}$  satisfy

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Optimal: replacing the inequality with an equality, this is exactly the condition leading to bifurcations.

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$\tilde{W}_{k^*} = \frac{\sigma_0}{\Phi'_0 M_{\infty}},$	
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Goal: estimates for a combination of the time derivatives of the quantities

$$\mathcal{E} = \int_0^{+\infty} \left(\frac{\rho - \rho_\infty}{\rho_\infty}\right)^2 \rho_\infty ds, \quad \text{and} \quad \mathcal{H} = (\Phi_{\bar{\rho}} - \Phi_0)(\bar{\rho} - \bar{\rho}_\infty),$$

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Why? Linearizing  $\Phi-\Phi_0$  , we get

$$\int_{\mathbb{T}^d} \mathcal{E} - \frac{\mathcal{H}}{\sigma} dx \simeq \int_{\mathbb{T}^d} \int_0^\infty \left(\frac{\rho - \rho_\infty}{\rho_\infty}\right)^2 \rho_\infty \, ds dx \\ - \frac{\Phi'_0}{\sigma} \int_{\mathbb{T}^d} W * (\bar{\rho} - \bar{\rho}_\infty)(\bar{\rho} - \bar{\rho}_\infty) dx.$$

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Let  $\Phi$  be  $C^2$ , and  $W \in L^2(\mathbb{T}^d)$  be component-wise symmetric. Assume that that for a suitably small  $\alpha > 0$ ,

$$\int_{\mathbb{T}^d} (1-\alpha)g^2(x) - \frac{\Phi_g^\delta(x)}{\sigma}g(x)\,dx > 0, \quad g \in L^2(\mathbb{T}^d),$$

where

$$\Phi_g^{\delta}(x) = \Phi \left( M_{\infty} W * g(x) + W_0 \bar{\rho}_{\infty} + B \right) - \Phi_0.$$

Then, the  $L^1$  norm of the relative entropy

$$\int_{\mathbb{T}^d} \int_0^{+\infty} \left( \frac{\rho(x,s,t) - \rho_{\infty}(s)}{\rho_{\infty}(s)} \right)^2 \rho_{\infty}(s) ds dx,$$

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$$\int_{\mathbb{T}^d} \int_0^{+\infty} \left( \frac{\rho(x,s,t) - \rho_{\rm c}}{\rho_{\infty}(s)} \right)^{\rm to the linear stability} \ {\rm condition \ modulo} \ \alpha:$$

decays exponentially fast whenever  $\rho_0$  relative entropy.

$$\hat{W}_k < \frac{\sigma(1-\alpha)}{M_\infty}$$

 $\Phi$  is linear it reduces

# Stability of spatially dependent stationary states? Let

$$M_{\infty}(x) = \int_0^{+\infty} (s - \bar{\rho}_{\infty}(x))^2 \rho_{\infty}(x, s) ds.$$

Assume that

$$\|\Phi'\|_{\infty}\|W\|_{L^{2}(\mathbb{T}^{d})}\sup_{x\in\mathbb{T}^{d}}M_{\infty}^{1/2}(x)<\frac{\sigma}{2}\tilde{\gamma}(\rho_{\infty})^{1/2},$$

where  $\tilde{\gamma}(\rho_{\infty}) = \inf_{x \in \mathbb{T}^d} \gamma(\rho_{\infty}(x))$ , and  $\gamma(\rho_{\infty}(x))$  is the Poincaré constant for  $\rho_{\infty}(x)$ . Then,

$$\int_{\mathbb{T}^d} \int_0^{+\infty} \left( \frac{\rho(x,s,t) - \rho_{\infty}(x,s)}{\rho_{\infty}(x,s)} \right)^2 \rho_{\infty}(x,s) ds dx,$$

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Assume that

$$\|\Phi'\|_{\infty}\|W\|_{L^{2}(\mathbb{T}^{d})} < \frac{1}{2},$$

Then,

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# Summary and outlook

We

- know there exists spatially heterogeneous stationary states,<sup>2</sup>
- have almost optimal conditions for (local) stability of the spatially homogeneous stationary state, and
- have established stability of (possibly spatially heterogeneous) stationary states under somewhat restrictive conditions,

but what about

- existence and stability of the hexagonal states, or
- multistability?



<sup>2</sup>Carrillo, Roux, S., Physica D 2023.

# Thank you!

A preprint with J. Carrillo and P. Roux will be on arxiv very soon. (: