## TFY4305 solutions exercise set 5 2014

## Problem 3.5.4

a) Looking at Fig. 1 in the textbook, we see that the restoring force is given

$$F = (\sqrt{h^2 + x^2} - L_0)k .$$
 (1)

The force is in the direction of the spring. We need the component of the force along the wire, which is given by

$$F_{\text{wire}} = F \sin \theta$$
  
=  $F \frac{x}{\sqrt{h^2 + x^2}}$ , (2)

Newton's second law of motion then becomes

$$m\ddot{x} + b\dot{x} + kx\left(1 - \frac{L_0}{\sqrt{h^2 + x^2}}\right) = 0.$$
 (3)

b) Equilibrium solutions are given by  $\ddot{x} = \dot{x} = 0$ . This yields

$$kx\left(1 - \frac{L_0}{\sqrt{h^2 + x^2}}\right) = 0.$$

$$\tag{4}$$

Thus x = 0 and  $1 - \frac{L_0}{\sqrt{h^2 + x^2}} = 0$ . This gives  $\underline{x = 0}$  and  $\underline{x = \pm \sqrt{L_0^2 - h^2}}$ . The nonzero solutions exist only if  $L_0 \ge h$ .

c) If m = 0, the equation of motion reads

$$\dot{x} = \frac{kx}{b} \left( \frac{L_0}{\sqrt{h^2 + x^2}} - 1 \right) . \tag{5}$$

The function that determines the stability of the fixed points is

$$f(x) = \frac{kx}{b} \left( \frac{L_0}{\sqrt{h^2 + x^2}} - 1 \right) .$$
 (6)

The derivative is given by

$$f'(x) = \frac{k}{b} \left( \frac{L_0}{\sqrt{h^2 + x^2}} - 1 - \frac{L_0 x^2}{(k^2 + x^2)^{\frac{3}{2}}} \right) .$$
(7)

This yields

$$f'(0) = \frac{k}{b} \left(\frac{L_0}{h} - 1\right) \tag{8}$$

The origin is stable for  $L_0 < h$  and unstable for  $L_0 > h$ . Moreover

$$f'(\pm\sqrt{L_0^2 - h^2}) = \frac{k}{b} \left(1 - \frac{L_0^2}{h^2}\right) .$$
(9)

Thus the fixed points  $x = \pm \sqrt{L_0^2 - h^2}$  are stable since  $L_0 \ge h$ . The stability of the fixed points depends on the ratio  $y \equiv L_0/h$ . We therefore plot the bifurcation diagram as a function of this dimensionles variable.



Figure 1: Bifurcation diagram for the bead on a wire. The origin changes stability in x = 0.

d) We can rewrite the equation of motion as

$$\frac{m}{k}\frac{d^2x}{dt^2} + \frac{b}{k}\frac{dx}{dt} + x\left(1 - \frac{L_0}{\sqrt{h^2 + x^2}}\right) = 0.$$
(10)

Defining a new dimensionless time variable  $\tau$  via

$$t = \frac{b}{k}\tau . (11)$$

In terms of the new time variable, Newton's equation reads

$$\frac{m}{k}\frac{k^2}{b^2}\frac{d^2x}{d\tau^2} + \frac{dx}{d\tau} + x\left(1 - \frac{L_0}{\sqrt{h^2 + x^2}}\right) = 0.$$
(12)

One can therefore ignore the second-order term if

$$\frac{\underline{mk}}{\underline{b^2}} \ll 1 . \tag{13}$$

## Problem 3.6.5

The forces acting on the bead are gravity and the force from the spring (see Fig. 3.6.7). The component of gravity along the wire is  $F = mg \sin \theta$ . The force from the spring is

$$F_{\rm spring} = k \left( \sqrt{x^2 + a^2} - L_0 \right) .$$
 (14)

The component of this force along the wire is

$$F_{\text{wire}} = -\frac{x}{\sqrt{x^2 + a^2}} F_{\text{spring}} .$$
(15)

The sum of F and  $F_{\text{wire}}$  is zero when the system is in equilibrium. This yields

$$mg\sin\theta - kx\left(1 - \frac{L_0}{\sqrt{x^2 + a^2}}\right) = 0.$$
<sup>(16)</sup>

b) Dividing Eq. (16) by kx, we can write

$$\frac{mg}{kx}\sin\theta = \left(1 - \frac{L_0}{a}\frac{1}{\sqrt{1 + (\frac{x}{a})^2}}\right).$$
(17)

Rearranging terms, we find

$$1 - \frac{mg}{ka} \frac{a}{x} \sin \theta = \frac{L_0}{a} \frac{1}{\sqrt{1 + (\frac{x}{a})^2}}.$$
 (18)

This is on the form

$$1 - \frac{h}{u} = \frac{R}{\sqrt{1 + u^2}},$$
 (19)

if we identify  $\underline{u = x/a}$ ,  $\underline{R = L_0/a}$ , and  $\underline{h = \frac{mg}{ka} \sin \theta}$ .

c) Note that the variable u can be positive and negative. Without loss of generality we can restrict h to positive values (negative values correspond to tilting the wire the other way) The function g(u) = 1 - h/u approaches zero as  $u \to \pm \infty$ . The maximum of the function  $h(u) = R/\sqrt{1+u^2}$  is R. For R < 1, it is then clear that the functions intersect at a single point and hence there is one fixed point. This is shown in Fig. 2

For R > 1, the function g(u) still intersects the function h(u) for positive values of u. However, there may be one or two intersections for negative values of u depending on h. In fact, as Figs. 3 and 4 suggest, a saddle-node bifurcation is taking place at a critical value of h. Since the derivatives of the functions g(u) and h(u) are equal at the bifurcation, h and u satisfy the two equations

$$1 - \frac{h}{u} = \frac{R}{\sqrt{1 + u^2}}, \qquad (20)$$

$$\frac{h}{u^2} = -\frac{Ru}{(1+u^2)^{\frac{3}{2}}}.$$
(21)



Figure 2: Graphical analysis.  $R = \frac{1}{2}$  and  $h = \frac{1}{2}$ .



Figure 3: Graphical analysis. R = 2 and h = 0.44.

d) The right-hand-side of Eq. (19) can be expanded in powers of u. This yields

$$1 - \frac{h}{u} \approx (r+1)(1 - \frac{1}{2}u^2),$$
 (22)

where r = R - 1. Ignoring the term  $\sim ru^2$  which is much smaller than the term  $\sim u^2$ , we obtain

$$-\frac{h}{u} \approx r - \frac{1}{2}u^2 , \qquad (23)$$

or

$$\frac{\frac{1}{2}u^3 - ru - h = 0}{2}.$$
(24)

e) We first define the functions  $f(u) = \frac{1}{2}u^3 - ru$  and g(u) = h(u). The bifurcation takes place when f(u) = g(u) and f'(u) = g'(u). The latter equation gives

$$\frac{3}{2}u^2 - r = 0. (25)$$

This yields  $\underline{R(u) = 1 + \frac{3}{2}u^2}$ . The equation g(u) = h(u) gives  $\underline{h(u) = \frac{1}{2}u^3 - ru = -u^3}$ 



Figure 4: Graphical analysis. R = 2 and h = 0.4.

f) The exact equations for the bifurcation are given by Eqs. (20) and (21). Solving the latter with respect to  $\frac{h}{u}$  and inserting it into the former, one finds

$$1 + \frac{Ru^2}{(1+u^2)^{\frac{3}{2}}} = \frac{R}{\sqrt{1+u^2}}.$$
 (26)

Solving this equation with respect to R gives

$$\underline{R(u) = (1+u^2)^{\frac{3}{2}}}.$$
(27)

Using Eq. (20), we find

$$\frac{h(u) = -\frac{Ru^3}{(1+u^2)^{\frac{3}{2}}} = -u^3}{(1+u^2)^{\frac{3}{2}}} = -u^3$$
(28)

Note that the exact Eqs. (27) and (28) reduce to the approximate solutions upon expanding R(u) to first nontrivial order in u for small u.

g) and h) were not a part of the exercise.