

# Global stabilization of periodic orbits in chaotic systems by using symbolic dynamics

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**Abstract:** In this report, a control method for the stabilization of periodic orbits for a class of discrete-time systems that are topologically conjugate to symbolic dynamics is proposed and applied to a population model in an ecosystem and the Smale horseshoe map. A periodic orbit is assigned as a target by giving a sequence in which symbols have periodicity. As a consequence, it is shown that any periodic orbits can be globally stabilized by using arbitrarily small control inputs. This work is the first attempt to systematically design a control system based on symbolic dynamics in the sense that one estimates the magnitude of control inputs and analyzes the Lyapunov stability.

Keywords: Nonlinear system control, Asymptotic stabilization, Tracking, Application of nonlinear analysis and design

## 1. INTRODUCTION

Chaos, signifying randomness and irregularity, is ubiquitous in nonlinear dynamical systems. The hallmark of chaos is sensitive dependence of the system's state on initial conditions. That is, a small error in the initial conditions can lead to a large error in the state of the system after a finite time interval. In many practical situations it is desirable if chaos can be avoided. The OGY-method (Ott et al. [1990]) was proposed as the first method controlling chaos in 1990, and since then, much related research has been carried out. The principal purpose of chaos control is stabilization of a periodic orbit embedded in an attractor.

Symbolic dynamics is introduced in order to characterize the orbit structure of a dynamical system via infinite sequences of "symbols" (Moser [1973], Wiggins [1991]). The study on symbolic dynamics has a long history. The first application was shown in Hadamard [1898]'s work of geodesics on surfaces of negative curvature. Birkhoff [1927] used symbolic dynamics in his studies of dynamical systems. Morse and Hedlund [1938] studied symbolic dynamics as an independent subject. Levinson [1949] applied it for the study of the forced van der Pol equation, and from his result, Smale [1963] introduced the well-known horseshoe mapping. In chaos engineering, symbolic dynamics is used for chaos communication (Hayes et al. [1993]) and the targeting problem<sup>1</sup> (Glenn and Hayes [1996], Corron and Pethel [2003]).

The purpose of our study is the global stabilization of a periodic orbit embedded in an attractor. To this end, first, a control law is designed in the sequence space such that

<sup>1</sup> The targeting problem is the problem of how to drive trajectories from initial states to the neighborhood of a target orbit.

the target periodic orbit becomes asymptotically stable. Next, the control law is transformed to the state space. By the proposed method, we design a one-dimensional control system for a population model in an ecosystem, and a two-dimensional control system with one input for the Smale horseshoe map. Our work is the first exposition that uses symbolic dynamics in order to systematically design control systems. The use of symbolic dynamics for design is effective in the sense that it is possible to globally stabilize any periodic orbit with arbitrarily small inputs by a uniform control law that does not switch from targeting to local stabilization, which is not an easy task with the conventional state space approach.

## 2. SYMBOLIC DYNAMICS

### 2.1 Symbolic dynamics

Let us consider a discrete-time dynamical system that has an invariant set  $X$ .

$$x_{n+1} = f(x_n), \quad x_n \in X. \quad (1)$$

Let  $S = \{0, 1, \dots, N\}$  be a set of symbols, and let  $(N+1)$  subsets  $X_k$  for  $k \in S$  be disjoint sets, the union of which is the invariant set  $X$ .

$$X = X_0 \cup X_1 \cup \dots \cup X_N, \quad X_i \cap X_j = \emptyset \quad (i \neq j).$$

Let the symbol  $s_i \in S$  be as follows<sup>2</sup>.

$$f^i(x) \in X_k \implies s_i = k.$$

Let also the set  $\Sigma$  be the infinite direct product of  $S$ ,  $\Sigma := \prod_{i=-\infty}^{\infty} S_i$  ( $S_i = S$ ), which is called the *sequence space*. We define a mapping  $\Psi : X \rightarrow \Sigma$  by

$$\Psi(x) = \dots s_{-2}s_{-1} \bullet s_0 s_1 s_2 \dots, \quad x \in f^{-i}(X_k),$$

<sup>2</sup>  $f^i(\cdot)$  means taking the composition of  $f$  with itself  $i$  times.

and describe this infinite sequence  $\Psi(x)$  as  $\rho$ . Here, we let the symbol with a decimal point on the left of it be the one corresponding to a current state. Furthermore, define a mapping  $\sigma : \Sigma \rightarrow \Sigma$  as follows.

$$\sigma(\cdots s_{-2}s_{-1}\bullet s_0s_1\cdots) = \cdots s_{-1}s_0\bullet s_1s_2\cdots.$$

This mapping is called the *shift*.

Denote the dynamics of the mapping  $f$  on its invariant set  $X$  as  $(X, f)$ , and the dynamics of the mapping  $\sigma$  on  $\Sigma$  as  $(\Sigma, \sigma)$ . When  $\Psi$  is a homeomorphic mapping and satisfies  $\sigma \circ \Psi = \Psi \circ f$ , the pairs  $(X, f)$  and  $(\Sigma, \sigma)$  are said to be topological conjugate, which is represented by the commutative diagram in Fig. 1. Then, the system  $(\Sigma, \sigma)$  is called *symbolic dynamics* for the system  $(X, f)$ .

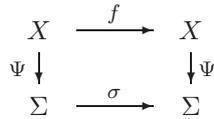


Fig. 1. The commutative diagram

*Remark 1.* The study on symbolic dynamics has a history of over a century. The advantages of using symbolic dynamics are as follows. If symbolic dynamics can be introduced for a dynamical system in the state space, the description of its time evolution in the sequence space, that is, shifting symbols, is simpler than that of the original system. It is easier to focus on certain properties of a dynamical system. For example, the existence of a periodic orbit with any period can be easily proven, and it is even possible to show there is a dense orbit in the state space.

*Remark 2.* The class of dynamics to which symbolic dynamics can be introduced is large. Actually, it is known that, for dynamics satisfying "the axiom A"<sup>3</sup>, the non-wondering set can be divided into finite basic sets and each of basic sets introduces Markov sub-shift<sup>4</sup> by "Markov partition" (See Robinson [1999]). Many dynamics, for example, Morse-Smale system, Anosov system, DA map and horseshoe, satisfy the axiom A.

### 2.2 Periodic orbits and the stability

*Definition 3.* If a trajectory  $\{x_0, x_1, \dots\}$  of the dynamics  $(X, f)$  satisfies that  $x_{n+T} = x_n$  for some constant  $T \in \mathbb{N}$ , the trajectory is called a *T-periodic orbit* or simply a *periodic orbit*. Then, each point of the T-periodic orbit is called a *T-periodic point* or a *periodic point*.

For a sequence in  $\Sigma$  corresponding to a periodic point in  $X$ , we have the following proposition.

*Proposition 4.* A state  $x \in X$  is a T-periodic point if and only if the sequence  $\rho \in \Sigma$  corresponding to  $x$  consists of infinitely repeated T-length blocks of symbols.

$$\rho = \Psi(x) = \cdots \underbrace{s_1s_2\cdots s_T}_{T\text{-length block}} \underbrace{s_1s_2\cdots s_T}_{T\text{-length block}} \cdots.$$

<sup>3</sup> A diffeomorphism  $f : M \rightarrow M$  ( $M$  is a manifold) which is such that the non-wondering set  $\Omega(f)$  is hyperbolic and a set of all periodic points  $Per(f)$  is dense in  $\Omega(f)$ , is said to satisfy axiom A.

<sup>4</sup> Markov sub-shift is a restriction of the shift  $\sigma$  to  $\Sigma_A$ , where  $\Sigma_A \subset \Sigma$  is a  $\sigma$ -invariant subset given by a transition matrix describing how sequences evolve.

From this proposition, it turns out that all periodic orbits in the invariant set can be specified by sequences.

In this report, we describe a periodic point in  $X$  and a sequence in  $\Sigma$  corresponding to it by adding "·", as  $\bar{x}$  and  $\bar{\rho} = \cdots \bar{s}_0\bar{s}_1\cdots := \Psi(\bar{x})$ , respectively. Furthermore, denote a T-Periodic orbit by a finite set  $\gamma_T = \{\bar{x}_0, \bar{x}_1, \dots, \bar{x}_{T-1}\}$ , and let  $\mathcal{P}_T$  be a set of the sequences corresponding to  $\gamma_T$  as follows.

$$\mathcal{P}_T = \{\bar{\rho}_0, \bar{\rho}_1, \dots, \bar{\rho}_{T-1}\}, \bar{\rho}_i = \Psi(\bar{x}_i), i = 1, 2, \dots, T.$$

We define the distance between a state  $x \in X$  and a periodic orbit  $\gamma_T \subset X$  by  $d(x, \gamma_T) := \min_{y \in \gamma_T} dist(x, y)$ , where  $dist$  is a metric in  $X$ . And also, we define the stability of a periodic orbit as follows.

*Definition 5.* A periodic orbit  $\gamma_T$  is said to be stable if, for all  $\varepsilon > 0$ , there exists a  $\delta = \delta(\varepsilon) > 0$  such that, for any solution  $\{f^n(x_0)\}$  satisfying  $d(x_0, \gamma_T) < \delta$ , we have  $d(f^n(x_0), \gamma_T) < \varepsilon$  for all  $n \geq 0$ . A periodic orbit is said to be unstable if it is not stable.

*Definition 6.* A periodic orbit  $\gamma_T$  is said to be globally asymptotically stable if it is stable and, for any initial state  $x_0$ , we have  $\lim_{n \rightarrow \infty} d(f^n(x_0), \gamma_T) = 0$ .

## 3. DESIGN OF A CONTROL SYSTEM BASED ON SYMBOLIC DYNAMICS

Now, let us consider the following control system for system (1).

$$x_{n+1} = f(x_n) + u_n. \tag{2}$$

We formulate the problem to be tackled in this report as follows.

*Problem* Design a control law (*i.e.*  $u_n$  in (2)) that globally stabilizes the unstable periodic orbit  $\gamma_T$  in system (1). Furthermore, design a control law that accomplishes the stabilization of  $\gamma_T$  with inputs whose magnitudes are less than a value given arbitrarily.

### 3.1 Control law in the sequence space

In the sequence space  $\Sigma$ , the time evolution of sequences by the shift mapping  $\sigma$  is described as

$$\rho_{n+1} = \sigma(\rho_n). \tag{3}$$

Here,  $\rho_n$  is the sequence corresponding to the state  $x_n$ .

Designing a control system that satisfies the requirements of the problem is equivalent to altering  $\sigma$  so that an orbit starting at an arbitrary initial sequence  $\rho_0$  converges to  $\mathcal{P}_T \subset \Sigma$  corresponding to  $\gamma_T \subset X$ . We notice, due to the metric<sup>5</sup> in  $\Sigma$ , that the more symbols from the decimal point toward both sides agree in  $\rho$  and  $\rho'$ , the closer  $\rho$  and  $\rho'$  are. Therefore, new  $\sigma$ , which we denote as  $\pi$ , requires rewriting symbols in the sequence. Let  $k$  and  $l$  be integers with  $k \geq 0, l \geq 1$ , respectively. Assume that each of the T-periodic sequences in  $\mathcal{P}_T$  consists of infinitely repeated

<sup>5</sup> The metric between two two-side infinite sequences  $\rho, \rho'$  is given by  $dist_\Sigma(\rho, \rho') := \sum_{i=-\infty}^{\infty} \frac{1}{2^{|i|}} \frac{|s_i - s'_i|}{1 + |s_i - s'_i|}$ .



Furthermore, since  $x_{n+1} = f(x_n) + u_n = f(\bar{x}_n)$ , it turns out that  $x_{n+1} = \bar{x}_{n+1} \in \gamma_T$ . ■

Furthermore, we have the proposition concerning the magnitude of the inputs of (6).

*Proposition 10.* For all  $\varepsilon > 0$ , there exists a  $K = K(\varepsilon) > 0$  such that, if the design parameter  $k$  is larger or equal to  $K$ , then we have  $\|u_n\| < \varepsilon$  for all  $n \geq 0$ .

**Proof.** The sequences corresponding to  $\tilde{f}(x_n)$  and  $f(x_n)$  are

$$\begin{aligned} \tilde{\rho}_{n+1} &:= \Psi(\tilde{f}(x_n)) = \cdots \underbrace{s_n}_{\parallel} \bullet \underbrace{s_{n+1}s_{n+2} \cdots s_{n+k}}_{\parallel} \bar{s}_{n+k+1} \cdots \\ \rho_{n+1} &:= \Psi(f(x_n)) = \cdots \underbrace{s_n}_{\parallel} \bullet \underbrace{s_{n+1}s_{n+2} \cdots s_{n+k}}_{\parallel} s_{n+k+1} \cdots, \end{aligned}$$

respectively, that is, all left symbols and at least  $k$  right symbols from the decimal point in two sequences agree. From the metric in  $\Sigma$ , it turns out that  $d_\Sigma(\tilde{\rho}_{n+1}, \rho_{n+1}) < 1/2^{k-1}$ . Since  $\Psi^{-1}$  is continuous, the smaller the distance between  $\tilde{\rho}_{n+1}$  and  $\rho_{n+1}$  is, the smaller the distance between  $\tilde{f}(x_n)$  and  $f(x_n)$  is. Therefore, given  $\varepsilon > 0$ , there exists a  $\delta$  such that, if  $d_\Sigma(\tilde{\rho}_{n+1}, \rho_{n+1}) < \delta$ , then  $dist(\tilde{f}(x_n), f(x_n)) < \varepsilon$ . If we choose  $k$  such that  $k > \log_2(1/\delta) + 1$ , then  $d_\Sigma(\tilde{\rho}_{n+1}, \rho_{n+1}) < \delta$ , and thus, we have  $\|u_n\| < \varepsilon$ . ■

#### 4.2 The stability analysis of periodic orbits

In order to analyze the stability of periodic orbits of the feedback system (6), we define the neighborhood  $V_j$  of a sequence  $\rho$  in  $\Sigma$  by

$$V_j(\rho) := \{\tilde{\rho} \in \Sigma \mid \tilde{s}_i = s_i, |i| < j\}.$$

For an integer  $j$  and a periodic orbit  $\gamma_T = \{\bar{x}_0, \bar{x}_1, \dots, \bar{x}_T\}$ , we define a maximum radius  $\varepsilon_j$  of a neighborhood of  $\gamma_T$  by

$$\varepsilon_j := \max_{0 \leq n \leq T-1} \sup_{\rho \in V_j(\bar{\rho}_n)} dist(\Psi^{-1}(\rho), \bar{x}_n),$$

where  $\bar{\rho}_n = \Psi(\bar{x}_n)$ . We have the following Lemma.

*Lemma 11.* For all integer  $j$ ,  $\varepsilon_j$  exists. If  $j' \geq j$ , then  $\varepsilon_{j'} \leq \varepsilon_j$ . Furthermore, we have  $\lim_{j \rightarrow \infty} \varepsilon_j = 0$ .

For the feedback system (6), we can prove the following proposition.

*Proposition 12.* Let  $l \geq 2$ . Then,  $\gamma_T$  is globally asymptotically stable.

**Proof.** From Proposition 9, it is proven that  $\gamma_T$  is a periodic orbit in the feedback system (6).

For given  $\varepsilon > 0$  and  $k \geq 0$ , let  $\delta = \min\{\varepsilon, 1/2^{k+1}\}$ . For  $\bar{\rho} = \cdots \bar{s}_0 \bar{s}_1 \cdots \in \mathcal{P}_T$ , if  $\rho = \cdots s_0 s_1 \cdots$  satisfies  $d_\Sigma(\rho, \bar{\rho}) < \delta$ , then we have

$$s_i = \bar{s}_i, |i| \leq \eta,$$

where  $\eta$  is the largest integer less than or equal to  $\max\{\log_2(1/\varepsilon) - 1, k\}$ . Since some symbols in  $\pi(\rho)$  and  $\pi(\bar{\rho})$  agree as follows,

$$\begin{aligned} \pi(\rho) &= \cdots \underbrace{\bar{s}_{-\eta} \cdots \bar{s}_0}_{\parallel} \bullet \underbrace{\bar{s}_1 \bar{s}_2 \cdots \bar{s}_\eta}_{\parallel} \underbrace{\bar{s}_{\eta+1} \cdots \bar{s}_*}_{\parallel} s_{*+1} \cdots \\ \pi(\bar{\rho}) &= \cdots \underbrace{\bar{s}_{-\eta} \cdots \bar{s}_0}_{\parallel} \bullet \underbrace{\bar{s}_1 \bar{s}_2 \cdots \bar{s}_\eta}_{\parallel} \underbrace{\bar{s}_{\eta+1} \cdots \bar{s}_*}_{\parallel} \bar{s}_{*+1} \cdots \\ &\quad \uparrow \qquad \qquad \qquad \uparrow \qquad \qquad \qquad \uparrow \\ &\quad (\eta+1)\text{-length} \quad \eta\text{-length} \quad \text{more than } l\text{-length} \end{aligned}$$

it turns out

$$d_\Sigma(\pi(\rho), \pi(\bar{\rho})) < d_\Sigma(\rho, \bar{\rho}) < \delta.$$

Therefore, we have

$$d_\Sigma(\pi^n(\rho), \pi^n(\bar{\rho})) < \delta \leq \varepsilon, \quad n \geq 0.$$

Note that, for a sequence  $\rho$ , if a periodic sequence  $\bar{\rho}_i$  is the closest to  $\rho$  in  $\mathcal{P}_T$ , then  $\pi(\bar{\rho}_i)$  is the closest to  $\pi(\rho)$  in  $\pi(\mathcal{P}_T)$ . Therefore, it turns out that, if  $d_\Sigma(\rho, \mathcal{P}_T) < \delta$ , then  $d_\Sigma(\pi^n(\rho), \mathcal{P}_T) < \varepsilon$  for all  $n \geq 0$ .

By the continuity of  $\Psi^{-1}$ , we prove that, for all  $\lambda > 0$ , there exists a  $\varepsilon = \varepsilon(\lambda) > 0$  such that, if  $d_\Sigma(\pi^n(\Psi(x)), \mathcal{P}_T) < \varepsilon$ , then  $d(\Psi^{-1}(\pi^n(\Psi(x))), \gamma_T) = d(\tilde{f}^n(x), \gamma_T) < \lambda$ . Similarly, by the continuity of  $\Psi$ , it turns out that, for all  $\delta > 0$ , there exists a  $\nu = \nu(\delta) > 0$  such that, if  $d(x, \gamma_T) < \nu$ , then  $d_\Sigma(\Psi(x), \mathcal{P}_T) < \delta$ . Therefore, one concludes that, for all  $\lambda > 0$ , there exists a  $\nu > 0$  such that, if  $d(x, \gamma_T) < \nu$ , then we have  $d(\tilde{f}^n(x), \gamma_T) < \lambda$ . The stability of  $\gamma_T$  is proven.

The global asymptotic stability is proven as follows. For an arbitrary initial state  $x_0$ , the sequence at time  $n(\geq k)$ ,  $\pi^n(\Psi(x_0))$ , has  $k + n(l - 1)$  ( $=: j_n$ ) symbols from the left being equal to those of a periodic sequence in  $\mathcal{P}_T$ . Therefore, a state  $x_n$  satisfies that  $d(x_n, \gamma_T) \leq \varepsilon_{j_n}$ . Since  $\lim_{n \rightarrow \infty} j_n = \infty$ , we have  $\lim_{n \rightarrow \infty} \varepsilon_{j_n} = 0$ . Therefore, it turns out that, for an arbitrary initial state  $x_0 \in X$ , we have  $\lim_{n \rightarrow \infty} d(x_n, \gamma_T) = 0$ . ■

*Remark 13.* The distance between  $x_n$  and  $\gamma_T$  converges to 0 more rapidly by choosing larger  $l$ .  $l \geq 2$  means, however, the calculation amount is getting larger actually. If one wants to avoid it, one has to let  $l = 1$ . We note that, although the asymptotic stability cannot be guaranteed for  $l = 1$ , the states can keep in a narrow tube around the periodic orbit by choosing reasonably large  $k$ .

## 5. APPLICATIONS

### 5.1 Control of an ecosystem

One of the simplest systems an ecologist can study is seasonally breeding populations in which generations do not overlap (May [1976]). For example, many natural populations such as temperate zone insects are of this kind. Such a relationship is expressed by a discrete-time system  $x_{n+1} = f(x_n)$  (variable  $x_n$  is the magnitude of the population). There are other examples expressed in this form, as, for example, in biology the theory of genetics and epidemiology. In economics the models for the relationship between commodity quantity and price and for the theory of business cycles. In sociology, the theory of learning and the propagation of rumors in variously structured societies are described by this kind of equation. In many of these contexts, and for biological populations in particular, there is a tendency for the variable  $x_n$  to increase from one generation to the next when it is small, and for it to

decrease when it is large. The discrete-time system below is a model representing such a tendency.

$$x_{n+1} = rx_n(1 - x_n), \quad x_n \in [0, 1] \quad (9)$$

This system is called *Logistic map*, and known to show chaotic behavior by choosing parameter  $r$  suitably. In particular, when  $r = 4$ , the system generates chaos (Robinson [1999]), and the closed interval  $[0, 1]$  is an invariant set. Furthermore, divide the interval  $[0, 1]$  into two regions with the boundary value  $1/2$  and give symbols "0" and "1" to the regions, respectively. That is, denote these regions as  $X_0 = [0, 1/2)$ ,  $X_1 = [1/2, 1]$ . Then, symbolic dynamics  $(\Sigma, \sigma)$  can be introduced into the system (9) with  $r = 4$ . Here, sequences are one-side infinite sequences because (9) is not invertible. However, since the control law rewrites symbols being located on right side from the decimal point, our control method can be applied to this system (9).

Now, by adding or removing individuals in (9), we try to fluctuate the population of the individuals periodically. In particular, it is intended that the magnitude of the population always returns to the initial magnitude every 3 generations. For such a purpose, we give a 3-periodic sequence repeating "011" as a target orbit and design a control system by using the proposed method. The simulation results are shown below. Fig. 4 illustrates the time evolution of the state starting at the initial condition  $x_0 = 0.3$  with no control. That is, a chaotic behavior can be observed. Fig. 5 and Fig. 6 show the time evolutions of the states starting at the same initial condition  $x_0 = 0.3$  with the design parameters  $(k, l) = (1, 2)$ ,  $(10, 2)$ , respectively. Also the state values (the magnitude of the population) are plotted in the top figures and the input values are plotted in the bottom figures, respectively. From Fig. 5 and 6, it is confirmed that the states converge to the 3-periodic orbit. Furthermore, by comparing Fig. 5 and 6, it can be verified that the system with the input magnitude parameter  $k = 10$  has smaller input values than those of the system with  $k = 1$ .

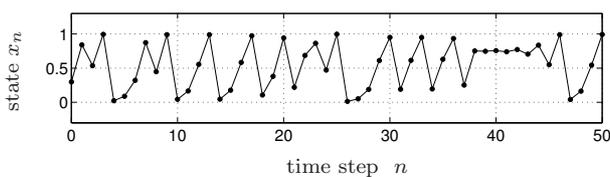


Fig. 4. Time evolution of the state without control input; initial condition  $x_0 = 0.3$

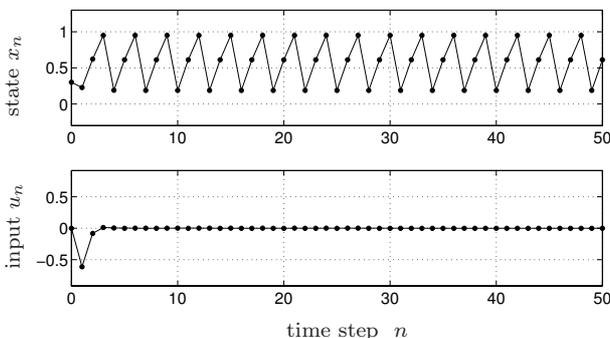


Fig. 5. Stabilization of a 3-periodic orbit embedded in a Logistic map; initial condition  $x_0 = 0.3$ ; design parameter  $k = 1, l = 2$

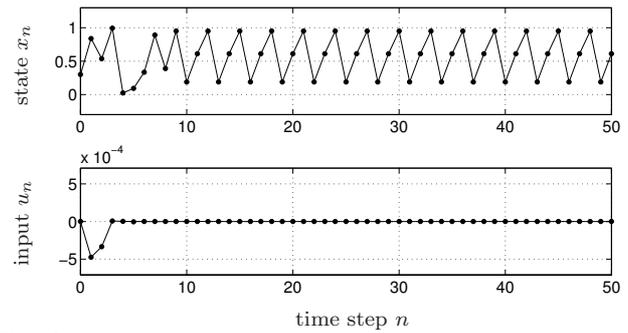


Fig. 6. Stabilization of a 3-periodic orbit embedded in a Logistic map; initial condition  $x_0 = 0.3$ ; design parameter  $k = 10, l = 2$

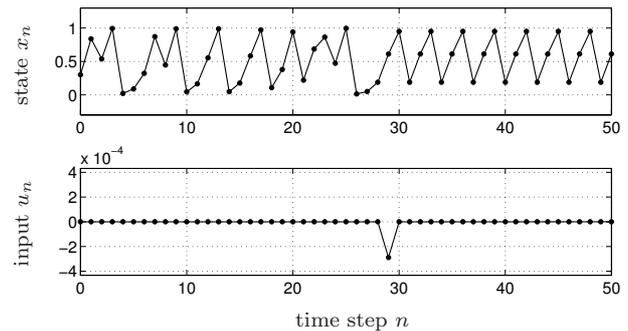


Fig. 7. Stabilization of a 3-periodic orbit embedded in Logistic map with the OGY-method; initial condition  $x_0 = 0.3$

*Comparison with the OGY-method* In the OGY-method (Ott et al. [1990]), the control inputs are added so that trajectories transit onto a local stable manifold, only when the state enters in a control region given in advance. The waiting time is derived statistically, and, in general, the smaller the control region is, the longer the time is. In our method, however, one can clearly know when the state will enter the neighborhood of the target even if it is small.

Fig. 7 illustrates a simulation result of stabilization of the 3-periodic orbit by applying the OGY-method. The control region is a neighborhood of the 3-periodic orbit with a radius 0.001. It can be verified that it takes longer time to stabilize the 3-periodic orbit than the proposed control method.

*A simulation of the feedback system with noise* For the logistic map (9), we consider a feedback system with white Gaussian noise  $\{v_n\}$  as follows.

$$x_{n+1} = f(x_n) + u(x_n) + v_n. \quad (10)$$

We set the mean and the standard deviation of noise  $\{v_n\}$  to 0 and  $10^{-4}$ , respectively, and simulate (10) in the case when (i)  $k = 5, l = 2$  and (ii)  $k = 10, l = 2$ . Fig. 8 shows the time evolutions of the states in these cases. From Fig. 8, it turns out that, the 3-periodic orbit is stabilized in the case (i), but it is not done in the case (ii). One concludes that, if the design parameter  $k$  is not sufficiently small, that is, the upper limit of the inputs is not sufficiently large, to remove the effect of the noise, then periodic orbits in (6) cannot be stabilized.

### 5.2 Control of the Smale horseshoe map—two-dimensional system with one input

Smale horseshoe map was introduced as the first example of diffeomorphisms that had an infinite number of peri-

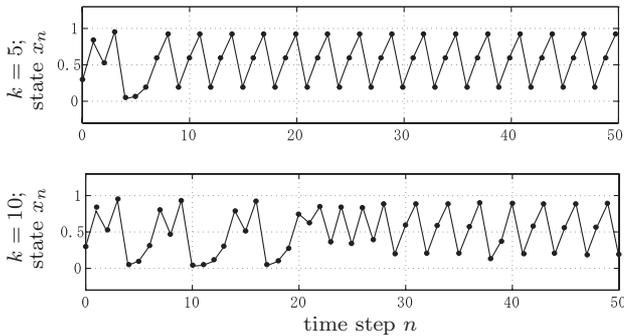


Fig. 8. Responses of the feedback system with white Gaussian noise: (a) Top figure; design parameter  $k = 5$ ,  $l = 2$ : (b) Bottom figure; design parameter  $k = 10$ ,  $l = 2$

odic points and were structurally stable (Smale [1963]). Furthermore, understanding of the Smale horseshoe is absolutely essential for understanding what is meant by term "chaos".

Consider a square  $D = [0, 1] \times [0, 1]$  on the plane, and the subsets  $H_0 = [0, 1] \times [0, 1/\mu]$  and  $H_1 = [0, 1] \times [1 - 1/\mu, 1]$ , where  $\mu > 2$ . The simplified Smale horseshoe map is given as follows (Wiggins [1991]).

$$\begin{pmatrix} x_{n+1} \\ y_{n+1} \end{pmatrix} = f(x_n, y_n), \quad (11)$$

$$f(x, y) = \begin{cases} \begin{pmatrix} \lambda & 0 \\ 0 & \mu \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}, & \begin{pmatrix} x \\ y \end{pmatrix} \in H_0 \\ \begin{pmatrix} -\lambda & 0 \\ 0 & -\mu \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} 1 \\ \mu \end{pmatrix}, & \begin{pmatrix} x \\ y \end{pmatrix} \in H_1, \end{cases} \quad (12)$$

where  $\lambda < 1/2$  and  $H_1, H_2$  are compressed in the direction of  $x$ -axis and stretched in the direction of  $y$ -direction. This system (11) has an invariant set  $\Lambda_{\lambda, \mu} = \{(x, y) \mid f^k(x, y) \in D, \forall k \in \mathbb{Z}\}$ , which is known to be a Cantor set. Let  $X_0 = H_0 \cap \Lambda_{\lambda, \mu}$  and  $X_1 = H_1 \cap \Lambda_{\lambda, \mu}$ . Then, symbolic dynamics can be introduced.

For the system (11) with  $\mu = 3$  and  $\lambda = 1/3$ , we try to stabilize a 4-periodic orbit in  $\Lambda_{1/3, 3}$  by the following two-dimensional control system with one input.

$$\begin{pmatrix} x_{n+1} \\ y_{n+1} \end{pmatrix} = f(x_n, y_n) + \begin{pmatrix} 0 \\ 1 \end{pmatrix} u_n, \quad (13)$$

where  $u_n$  is a scalar function. We give a 4-periodic sequence repeating "0011" as a target orbit. Fig. 9 shows the time evolutions of the states and the inputs with the initial value  $(x_0, y_0) = (1/9, 1/9) \in \Lambda_{1/3, 3}$  and the design parameter  $(k, l) = (5, 2)$ . From Fig. 9, it is confirmed that states converge to 4-periodic orbit.

## 6. CONCLUSION

In this report, for a class of discrete-time systems that are topologically conjugate to symbolic dynamics, we proposed a control method to stabilize periodic orbits. We also showed application examples of the proposed control method, which were a one-dimensional control system for a population dynamics represented by a Logistic map and a two-dimensional control system with one input for the Smale horseshoe map. This is the first attempt to design control systems by using symbolic dynamics systematically in the sense that one estimates the magnitude of control inputs and analyzes the Lyapunov stability. The proposed

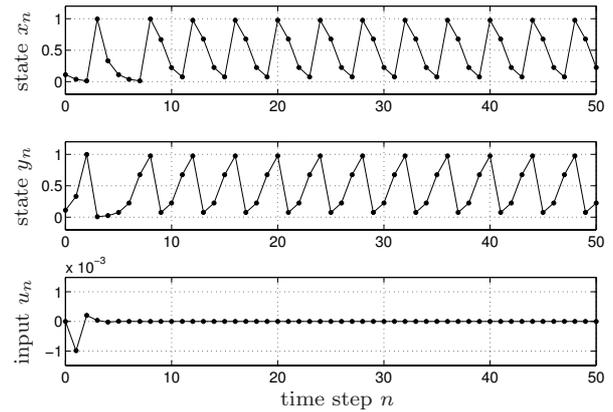


Fig. 9. Stabilization of a 4-periodic orbit embedded in a Smale horseshoe map; initial condition  $(x_0, y_0) = (1/9, 1/9)$ ; design parameter  $k = 5$ ,  $l = 2$

control method can stabilize any periodic orbits with arbitrarily small inputs without switching the control law from targeting to local stabilization, and can ensure the robustness against noise by choosing the design parameter suitably. It is difficult, with the conventional state space approaches, to accomplish the stabilization like this, showing the effectiveness of the use of symbolic dynamics.

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