Stress-Induced Crystallization of Poly(Trimethylene Terephthalate) Fibers by Molecular Dynamic Simulations

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Introduction

The phenomena of stress-induced (or flow-induced, strain-induced, force-induced) crystallization is characterized by the preferential orientation of polymer segments along the direction of applied force. Many researches report that the drawing and spinning process would increase the crystallinity of polymeric materials, resulting in oriented segments that would explosively improve the crystallization by oriented nuclei.

There are many works showing the evidences of the structural changes during drawing and spinning processes. From the behavior of stress-draw ratio curve of the melt-quenched amorphous PTT film during uniaxial drawing, a significant strain hardening could be observed simultaneously with crystallinity (measured by density) abruptly rising after draw ratio of 2.5.\textsuperscript{1,2} Besides, the development of oriented structure has been quantified by orientation factor calculated from IR dichroic ratio of characteristic bands\textsuperscript{3} and intensity of WXR\textsuperscript{4} of the melt-quenched amorphous PTT. Moreover, the birefringence of PTT increases with spinning speeds.\textsuperscript{5} On the other hand, many theoretical works on stress-induced crystallization using molecular simulations, most of which focus on linear polyethylene (PE). It is found that the torsional ordering\textsuperscript{6} and is followed by increase in orientation factor\textsuperscript{7} in the simulation with the pre-oriented models. The shish-kebab crystallization was observed by extensional flow process,\textsuperscript{8} and the secondary nucleation around a knot in the polymer system has been studied.\textsuperscript{9}

The change in the torsional conformation of a propylene glycol segment (O-CH\textsubscript{2}-CH\textsubscript{2}-CH\textsubscript{2}-O) during the drawing or spinning process is of great interest. The torsion angle distribution of the segment has been reported as \textit{tans-gauche-gauche-trans} (t-g-g-t) in the crystalline state.\textsuperscript{10,11} However, what the distribution in amorphous state is and how they transfer form amorphous to crystalline state remain unclear. The conformation of amorphous state may either be dominated by \textit{g-t-t-g}, state\textsuperscript{3} or in random\textsuperscript{12} states; some even estimated the content of the gauche state of $\phi_2$ to be 33.4\%.\textsuperscript{2} Furthermore, the mechanism of the conformation transition from amorphous to crystalline phase has been studied using IR spectroscopy\textsuperscript{2,3,12-14}, X-ray scattering\textsuperscript{15}, and $^{13}$C solid-state NMR.\textsuperscript{16} From those results, in crystalline phase, the $\phi_2$ exists largely in the \textit{gauche} state, and $\phi_1$ conformation exists largely in the \textit{trans} state. In amorphous phase, $\phi_2$ conformation exists mostly in the \textit{trans} state, but $\phi_1$ conformation is unclear. Moreover, during the annealing process, $\phi_2$ conformation changes form \textit{trans} state to \textit{gauche} state but $\phi_1$ conformation is unclear. During drawing process, $\phi_2$ conformation changes form \textit{trans} state to \textit{gauche} state, if polymer chains undergo sufficient thermal relaxation; however, $\phi_1$ conformation is unclear.

In this work we focus on the development of ordering of structures in the stress-induced crystallization process. (The thermal effect of the development ordering structure has been studied in our previous works.) The temperature dependence of the amount of ordered clusters, the behavior of ordered structures and torsion angle transition during drawing and annealing processes are analyzed. The results
presented here help us better understand the mechanism of the nucleation at the early stage of crystallization.

**Computational methods**

Atomistic molecular dynamic simulations are used to study the dynamic behavior of the nucleation process. The molecular models of PTT are prepared using commercial package Cerius2. Each model contains 4 chains of PTT molecules, each having a degree of polymerization of 27. Thus, each unit cell contains 108 repeating units of PTT, or equivalently 2708 atoms. In all the simulations, the generic force field Dreiding is used to describe the interactions among atoms in a system. The potential energy includes the bond stretching energy (U\(_b\)), angle bending energy (U\(_\theta\)), torsional angle rotation energy (U\(_\phi\)), electrostatic interaction (U\(_{\text{coul}}\)), and van der Waals interactions (U\(_{\text{vdw}}\)). Atomic (Mulliken) charges are determined from density functional theory (DFT) calculations for the monomer of PTT using the B3LYP functional and 6-31G** basis set in Gaussian 98. Moreover, the torsion terms of the propylene glycol segment (O-CH\(_2\)-CH\(_2\)-CH\(_2\)-O) being modified to better reproduce quantum mechanics (QM) results, according the energy barrier results of rotating torsion angles. This force field could reproduce reasonably the thermal properties (T\(_m\) and T\(_g\)), mechanical property (Young’s modulus), and crystal structural properties.

The computer code LAMMPS is used for all subsequent molecular dynamic simulations. A series of expansion and compression steps were applied to the initial structures in order to obtain equilibrium samples at 600K (which is above the melting temperature T\(_m\)). The 100 ps equilibrated sample is then cooled to 50 K at a rate of 1 K/ps. At each 50 K interval during the quenching process, the samples are then drawn with different drawing speeds (1x10\(^8\) s\(^{-1}\), 1x10\(^9\) s\(^{-1}\), 1x10\(^10\) s\(^{-1}\), and 1x10\(^11\) s\(^{-1}\)), and isothermally relaxed at several draw ratios (DR=1, 2, 3, and 4) up to 12 ns with pressure at 0 bar and N\(_T\)L\(_x\)σ\(_{yy}\)σ\(_{zz}\) ensemble with zero normal stress in y and z directions. Nose-Hoover thermostat is used for temperature control. The integration time step used is 0.5 to 1 fs.

**Results and discussions**

**Bulk Properties**

The changes in bulk properties are observed (Fig. 1) during stress-induced crystallization in our simulation. The increase of the amount of precursors in the draw process is much more remarkable than that in isothermal crystallization. Furthermore, the subsequent process of relaxation (NVT simulation after the drawing process) would lead to further increase in the fraction of oriented precursors. The rapid density decrease in each simulation indicates that the defects (voids) in the bulk phase are created during the drawing processes. At high draw speeds (1x10\(^8\) s\(^{-1}\), and 1x10\(^10\) s\(^{-1}\)) and the temperatures below T\(_g\) the formation of defects is prominent; on the other hand, at the slower draw speeds (1x10\(^8\) s\(^{-1}\), and 1x10\(^9\) s\(^{-1}\)) and temperatures above T\(_g\) the system relaxation is fast enough to reduce the amount of defects.

The changes in van der Waals interaction and the torsional energy reflect the transition of the structure under various drawing conditions. The van der Waals interaction increases with DR in all cases. These interaction energies also have a significant effect in the system density. For example, the
relaxation of van der Waals interaction leads to a constant density at lower draw speed ($1 \times 10^9 \text{ s}^{-1}$). A rapid increase in van der Waals interaction leads to a rapid drop in density at higher draw speed (1x10$^{11}$ s$^{-1}$). Besides, the van der Waals interaction increases with temperature at constant DR. The enhancement of van der Waals interactions slows down at temperatures higher than $T_g$ at lower draw speed. These results indicate that the mechanism of the chain packing is different in the various draw speeds and temperatures. This could be understood as a competition between thermal effect, which is prominent at slower draw speed and higher temperatures, and stress effect, which is prominent at faster draw speed and lower temperatures. On the other hand, the torsional energy decreases with DR at lower draw speeds ($1 \times 10^9 \text{ s}^{-1}$, and $1 \times 10^{10} \text{ s}^{-1}$), but fluctuates at the highest draw speed (1x10$^{11}$ s$^{-1}$). It is almost a linear function of temperature at constant DR. These results support that the stress effect influences the backbone chain rotation much more than thermal effect.

**Figure 1** The energy and density of the overall system.
The Fraction of Ordering Structure

The fraction of ordering structure, $X$, has been defined in previous work.\textsuperscript{20} As shown in Fig. 2, at small DR (DR=2), the fraction of precursors increases in all temperatures with slower draw speeds ($1\times10^9\text{s}^{-1}$ and $1\times10^{10}\text{s}^{-1}$), but decreases at temperatures 350K and 400K with higher draw speed ($1\times10^{11}\text{s}^{-1}$). At larger DR (DR>3), the saturation of induced precursor is observed for all temperatures with the slowest draw speed ($1\times10^9\text{s}^{-1}$). On the other hand, the development of precursors continues at the higher draw speed ($1\times10^{10}\text{s}^{-1}$ and $1\times10^{11}\text{s}^{-1}$). Besides, the orientation factor increases with DR, and the curves are similar in all situations, indicating that the orientation factor may be a function of draw ratio only. This implies that the stress-induced precursor is similar at the same DR with each draw speed.

The subsequent thermal relaxation (following the drawing process) could stabilize and facilitate in the growth of oriented precursors. The results of the fraction of precursors shows a temperature dependence similar to that of the isothermal crystallization, but the position of highest amount of precursors shifts to a lower temperature (300K with drawing vs 400 K without drawing\textsuperscript{20}). The local motion of polymer segments would be induced by uni-axis stress. (Note the maximum capacity of precursor formation should be found between the temperature allowing for short rage motion and that for entire chain motion.)

Figure 2 The fraction of ordering structure was shown as the function of temperature at a) $1\times10^9\text{s}^{-1}$ and b) $1\times10^{11}\text{s}^{-1}$, and c) the function of draw ratio at 400K with all draw speeds.
The torsional angle distribution

To understand the distribution of torsional angles in the PTT backbone (O-CH\(_2\)-CH\(_2\)-CH\(_2\)-O), we analyze the percentage of the trans and gauche states during the stress-induced crystallization. Different behaviors observed at different draw speeds. The percentage of \( \phi_1 \) and \( \phi_2 \) in trans state fluctuates at small DR, and increases at larger DR with slower draw speed. At higher draw speeds, this percentage decreases at small DR and increases at larger DR. The percentage value rebounds to the original value with draw speed 1x10\(^{11}\) s\(^{-1}\). On the other hand, the percentage of gauche state decreases with DR at all draw speeds. During thermal relaxation, the percentage of the trans rebounds in all draw speeds, but that of gauche rebounds only in faster draw speed. These results suggest that torsional angles change to transition state in the beginning drawing process, and then become trans state at higher DR or under thermal relaxation.

Detailed analysis also is performed on the two feature torsions (\( \phi_1 \) and \( \phi_2 \)) in the PTT backbone (Fig. 3). It is found that the trans-to-gauche ratio in \( \phi_1 \) increases rapidly (from 3 to 6) at lower drawing speed and higher temperatures; however, the increase is less significant at high drawing speed and low temperatures. In contrast, the gauche-to-trans ratio for \( \phi_2 \) seems to be insensitive to processing conditions (varies from 1.6 to 0.8 regardless of drawing speeds). After thermal relaxation, the value of trans-to-gauche ratio in \( \phi_1 \) improves significantly in the slower draw speeds, but remains almost constant in the fastest draw speed. On the other hand, the gauche-to-trans ratio for \( \phi_2 \) decreases at small DR with slower draw speeds (1x10\(^9\) s\(^{-1}\) and 1x10\(^{10}\) s\(^{-1}\)), but increases at larger DR with faster draw speed.

Before drawing, the fraction of trans state is ~34%, and that of gauche is ~25% in the bulk phase; the most populated conformations in \( \phi_1 \) is trans, and that in \( \phi_2 \) is gauche; i.e., the most populated conformations in the backbone torsions are t-g-g-t (\( \phi_1-\phi_2-\phi_2-\phi_1 \)). During draw process, the t-t-t-t conformation increases. Our results are in agreement with the result of Chuch’s experimental observation.\(^{14}\) The thermal relaxation would change the conformations in some of our simulations, but the changes are not obvious.

The states of backbone torsional angles in precursors evolve with DR differently in different processing conditions. The fraction of \( \phi_1 \) in trans increases with DR, but the fraction of \( \phi_2 \) in gauche decreases with DR in most precursors. The stress and thermal effects are different in these two torsional angles. Slower draw speed and lower temperature would improve the change of \( \phi_1 \) to the trans state; on the other hand, the faster draw speed and medium temperatures would improve the change of \( \phi_2 \) to the gauche state. Furthermore, the fraction of \( \phi_1 \) in trans is very high (~70%) in each precursor, but the fraction of \( \phi_2 \) in gauche only 20~30%. That indicates that the stress-induced transition in \( \phi_2 \) is faster than that in \( \phi_1 \). In addition, the faction of \( \phi_2 \) in gauche fluctuates in some cases, indicating the torsional angle \( \phi_2 \) would favorably move to the gauche after forcing to the trans state by the external stress.
Conclusion

The formation of oriented precursor in the stress-induced crystallization is observed in our simulation. The torsional and van der Waals forces are the dominant interactions, similar to the case observed in isothermal crystallization. The amount of stress-induced precursor increases in all regions of temperature. The maximum size of oriented precursor is larger than that created only by thermal stimulation. During stress-induced crystallization, the torsional distribution of the polymer backbone for segments rapidly rearrange to the t-t-t-t conformation in bulk phase. Within oriented precursors, the response of the torsional angle induced by stress is faster than that only induced by thermal stimulation, especially trans in $\phi_1$ (the transition rate of trans state in $\phi_1$ is faster than that of gauche state in $\phi_2$).

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References

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