

## TFY4205 Quantum Mechanics II

NTNU

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## SUGGESTED SOLUTION

**Problem 1**

The deviation takes the form

$$\lambda H_1 = \begin{cases} \frac{Ze^2}{4\pi\epsilon_0} \left( \frac{1}{r} - \frac{3}{2R} + \frac{r^2}{2R^3} \right) & \text{for } r \leq R \\ 0 & \text{for } r > R. \end{cases} \quad (1)$$

The contribution in first order perturbation theory is the same as the expectation value of the perturbation in the unperturbed state:

$$\langle 0 | \lambda H_1 | 0 \rangle = \int \lambda H_1(r) (\psi_{100})^2 d^3r. \quad (2)$$

The integrand is spherically symmetric and using standard techniques we arrive at

$$\langle 0 | \lambda H_1 | 0 \rangle = \frac{R^2}{a^2} \frac{Ze^2}{10\pi\epsilon_0 a}. \quad (3)$$

The total result for the energy eigenvalue is then

$$E_1 = E_1^0 \left( 1 - \frac{4}{5} \frac{R^2}{a^2} \right). \quad (4)$$

Since  $R/a \ll 1$ , this correction is very small (magnitude  $10^{-9}$  eV). Since  $E_1^0$  is negative, the correction will increase the ground state energy. This is reasonable since the perturbation increases the Coulomb potential for  $r < R$ . For larger  $Z$ , the correction will be more important due to the finite size of the nucleus, since  $E_1^0 \propto Z^2$ ,  $a^{-2} = Z^2/a_0^2$ , and because  $R$  is larger. Again, this is physically reasonable since the larger the nuclear charge, the closer the electron will be attracted to the nucleus, where it feels the deviation from a pure Coulomb potential.

**Problem 2**

According to the variational method, an upper estimate for the ground state energy is

$$E[f] = \frac{\int_{-\infty}^{\infty} f^*(x) \hat{H} f(x) dx}{\int_{-\infty}^{\infty} f^*(x) f(x) dx}. \quad (5)$$

Since our trial functions satisfy  $f(\pm\infty) = 0$ , we have

$$-\frac{\hbar^2}{2m} \int_{-\infty}^{\infty} f^* \frac{d^2 f}{dx^2} dx = \frac{\hbar^2}{2m} \int_{-\infty}^{\infty} \frac{df^*}{dx} \frac{df}{dx} dx. \quad (6)$$

The trial functions are real, meaning that we can compute

$$E[f] = \frac{\int_{-\infty}^{\infty} [\frac{\hbar^2}{2m} (f')^2 + V(x)f^2] dx}{\int_{-\infty}^{\infty} f^2 dx}. \quad (7)$$

For  $f(x) = e^{-\lambda x^2}$ , we find that

$$E[f] = \frac{\frac{\hbar^2}{2m} \sqrt{\frac{\pi\lambda}{2}} - \alpha}{\sqrt{\frac{\pi}{2\lambda}}}. \quad (8)$$

The minimum of this functional can be found by differentiating with respect to  $\lambda$  or by rewriting the expression. Using the latter procedure, we see that  $E[f]$  can be written as

$$E[f] = \left( \sqrt{\frac{\hbar^2\lambda}{2m}} - \alpha \sqrt{\frac{m}{\pi\hbar^2}} \right)^2 - \frac{m\alpha^2}{\pi\hbar^2}. \quad (9)$$

Only the first term is dependent on  $\lambda$ , so the minimum is obtained when this quadratic term vanishes. This determines the optimal value for  $\lambda$ , and thus the best estimate for the ground state energy is

$$E = -\frac{m\alpha^2}{\pi\hbar^2}. \quad (10)$$

For the trial function  $f(x) = e^{-\lambda|x|}$ , a similar calculation yields

$$E[f] = \frac{\frac{\hbar^2}{2m}\lambda - \alpha}{1/\lambda} = \left( \frac{\hbar\lambda}{\sqrt{2m}} - \frac{\sqrt{2m}\alpha}{2\hbar} \right)^2 - \frac{m\alpha^2}{2\hbar^2}. \quad (11)$$

Again, we can choose  $\lambda$  so that the quadratic term vanishes. Since  $-\frac{m\alpha^2}{2\hbar^2} < -\frac{m\alpha^2}{\pi\hbar^2}$ , the estimate obtained using  $f(x) = e^{-\lambda|x|}$  is the best one.

### Problem 3

To solve this problem, we must calculate the expectation value of the Hamiltonian for  $z \geq 0$ :

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dz^2} + Fz \quad (12)$$

using the trial function  $f(z)$ :

$$E[f] = \frac{\int_0^{\infty} f^* \hat{H} f dz}{\int_0^{\infty} |f|^2 dz}. \quad (13)$$

We will need three integrals which can be solved by standard techniques:

$$\begin{aligned} \int_0^{\infty} |f|^2 dz &= \alpha^{-3/2} \sqrt{\pi}/4, \\ -(\hbar^2/2m) \int_0^{\infty} f f'' dz &= (\hbar^2/2m) \frac{3}{8} \sqrt{\pi/\alpha}, \\ F \int_0^{\infty} z |f|^2 dz &= \frac{1}{2} F \alpha^{-2}. \end{aligned} \quad (14)$$

This gives us

$$E[f] = \frac{\hbar^2}{2m} \frac{3\alpha}{2} + \frac{2F}{\sqrt{\pi\alpha}}. \quad (15)$$

To find the minimum, we calculate the derivative with respect to  $\alpha$  and set it to zero. This gives:

$$\alpha = \left( \frac{2}{3\sqrt{\pi}} F \frac{2m}{\hbar^2} \right)^{2/3}. \quad (16)$$

Substituting this back into  $E[f]$ , we find the minimum value

$$E = 3 \left( \frac{3}{2\pi} \right)^{1/3} \left( \frac{\hbar^2}{2m} \right)^{1/3} F^{2/3}. \quad (17)$$